

$$q_n(t) = A e^{i\omega_n t} \Rightarrow a_n \cos(\omega_n t) + b_n \sin(\omega_n t)$$

Formulazione Lagrangiana per sistemi continui

Un sistema elastico di equazione

$$\frac{\partial^2 u}{\partial t^2} - \frac{\partial^2 u}{\partial x^2} = 0$$

è associata ad una Lagrangiana integrale

$$L = \frac{1}{2} \int \left[\left(\frac{du}{dt} \right)^2 - \left(\frac{du}{dx} \right)^2 \right] dx$$

Vestiamo come costituisce un metodo razionale che
riguarda direttamente la soluzione dell'eq. evoluzione.

Consideriamo L , la lagrangiana di un sistema continuo
come un funzionale

$$L(u, \frac{\partial u}{\partial x}, \frac{\partial u}{\partial t}, x, t)$$

Potremmo così vedere la lagrangiana totale di un
sistema continuo

$$L = \iint L(u, \frac{\partial u}{\partial x}, \frac{\partial u}{\partial t}, x, t) dx dt$$

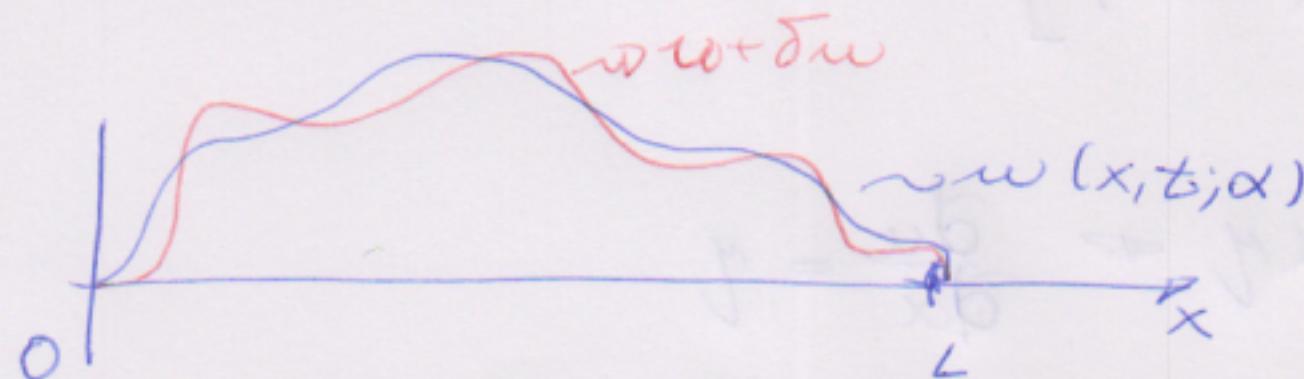
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dunque lagrangiana

In analogia col princ. variaz. di Ham., consideriamo le variazioni di al variare dell'incognita w e delle sue derivate)

Nota: adoro le variabili spazio e tempo non vengono variate

Sist. discreto \Rightarrow Variaz. di traiettorie del moto

Sist. continuo \Rightarrow Variaz. delle funzioni incognite



$$w(x,t;\alpha) = w(x,t,0) + \underbrace{\alpha \gamma(x,t)}_{\text{La variazione}}$$

in analogia al caso discreto al bordo $\gamma(0) = \gamma(L) = 0$: le variazioni non cambiano la c.c. del problema.

Definiamo

$$I(\alpha) = \iint dx dt L(w+\alpha\gamma, \partial_x w\gamma, \partial_t w\gamma, x, t) \alpha$$

Variazione di $I \Rightarrow \frac{dI}{d\alpha}$

Come nel caso elastico volevamo $\frac{dI}{d\alpha} \Big|_{\alpha=0}$

$$\frac{dI}{d\alpha} \Big|_{\alpha=0} = \delta \iint dx dt$$

$$\begin{aligned} \frac{dI}{d\alpha} &= \int_{t_1}^{t_2} \int_{x_1}^{x_2} dx dt \left[\frac{\partial L}{\partial u} \frac{du}{d\alpha} + \frac{\partial L}{\partial (\frac{\partial u}{\partial x})} \frac{d}{dx} \left(\frac{\partial u}{\partial x} \right) + \right. \\ &\quad \left. + \frac{\partial L}{\partial (\frac{\partial u}{\partial t})} \frac{d}{dt} \left(\frac{\partial u}{\partial t} \right) \right] \end{aligned}$$

$$u = u_0 + \alpha \gamma \Rightarrow \frac{du}{d\alpha} = \gamma$$

$$\frac{d}{dx} \left(\frac{\partial u}{\partial x} \right) = \frac{d}{dx} \left(\frac{\partial u_0 + \alpha \gamma}{\partial x} \right) = \frac{\partial \gamma}{\partial x}$$

$$\frac{d}{dt} \left(\frac{\partial u}{\partial t} \right) = \frac{\partial \gamma}{\partial t}$$

$$\frac{dI}{d\alpha} = \iint_{t_1 x_1}^{t_2 x_2} \left[\frac{\partial L}{\partial u} \cdot \gamma + \frac{\partial L}{\partial (\frac{\partial u}{\partial x})} \left(\frac{\partial \gamma}{\partial x} + \frac{\partial L}{\partial (\frac{\partial u}{\partial t})} \frac{\partial \gamma}{\partial t} \right) \right] dx dt$$

Integriiamo per parti

$$\iint_{t_1 x_1}^{t_2 x_2} \frac{\partial L}{\partial (\frac{\partial u}{\partial x})} \frac{\partial \gamma}{\partial x} dx dt = \int_{t_1}^{t_2} \left(\iint_{x_1}^{x_2} F(x, t) \frac{\partial \gamma}{\partial x} dx \right) dt$$

$$F(x, t) = \left. \frac{\partial L}{\partial \dot{u}} \right|_{\dot{u} = \dot{u}(x, t)}$$

$$\int_{x_1}^{x_2} F(x, t) \frac{\partial \eta}{\partial x} dx = \underbrace{\eta(x_1, t) F(x_1, t)}_{= 0} - \int_{x_1}^{x_2} \frac{\partial F}{\partial x} \eta dx$$

$$\Rightarrow \iint_{t_1, x_1}^{t_2, x_2} \frac{\partial L}{\partial \dot{u}} \frac{\partial \eta}{\partial x} dx dt = - \iint_{t_1, x_1}^{t_2, x_2} \frac{\partial}{\partial x} \left(\frac{\partial L}{\partial \dot{u}} \right) \eta dx dt$$

Analogamente

$$\iint_{t_1, x_1}^{t_2, x_2} \frac{\partial L}{\partial \dot{u}} \frac{\partial \eta}{\partial t} dx dt = - \iint_{t_1, x_1}^{t_2, x_2} \frac{\partial}{\partial t} \left(\frac{\partial L}{\partial \dot{u}} \right) \eta dx dt$$

Abbiamo ottenuto

$$\delta \iint_{t_1, x_1}^{t_2, x_2} L dx dt = \iint_{t_1, x_1}^{t_2, x_2} \left[\frac{\partial L}{\partial w} + \frac{\partial}{\partial x} \frac{\partial L}{\partial (\dot{u})} - \frac{\partial}{\partial t} \frac{\partial L}{\partial (\dot{u})} \right] \eta dx dt$$

$\delta \int L = 0 \forall \eta \Rightarrow$ Eq. EULERO LAGR per

sistemi conservativi

$$\boxed{\frac{\partial}{\partial t} \left(\frac{\partial L}{\partial \dot{u}} \right) + \frac{\partial}{\partial x} \frac{\partial L}{\partial (\dot{u})} - \frac{\partial L}{\partial w} = 0}$$

Analogia con E-L per sist. circolari

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{q}_i} - \frac{\partial L}{\partial q_i} = 0$$

Verif. che il principio Ham. fornisce l'eq. alle aule delle lag. ottenute con la teoria del oscillatori armati

$$L = \frac{1}{2} \iint \left[\left(\frac{\partial u}{\partial t} \right)^2 - \left(\frac{\partial u}{\partial x} \right)^2 \right] dx dt$$
$$\mathcal{L} = \frac{1}{2} \left[\left(\frac{\partial u}{\partial t} \right)^2 - \left(\frac{\partial u}{\partial x} \right)^2 \right]$$

$$\frac{\partial L}{\partial \left(\frac{\partial u}{\partial x} \right)} = -\frac{1}{2} \lambda \cdot \frac{\partial u}{\partial x} \Rightarrow \frac{\partial}{\partial x} \frac{\partial L}{\partial \left(\frac{\partial u}{\partial x} \right)} = -\frac{\partial^2 u}{\partial x^2}$$

$$\frac{\partial L}{\partial \left(\frac{\partial u}{\partial t} \right)} = \frac{1}{2} \lambda \cdot \frac{\partial u}{\partial t} \xrightarrow{\frac{\partial}{\partial t}} \frac{\partial^2 u}{\partial t^2}$$

$$\frac{\partial L}{\partial u} = 0$$

$$\Rightarrow -\frac{\partial^2 u}{\partial x^2} + \frac{\partial u}{\partial t^2} = 0$$

Notiamo che nel pr. Ham. le variabili x e t sono
sostituite in misura e con numeri se non sono

$$x_0 = t \quad x_1 = x$$

$$\Rightarrow \sum_i \frac{\partial}{\partial x_i} \frac{\partial L}{\partial (\frac{\partial u}{\partial x_i})} - \frac{\partial L}{\partial u} = 0$$

Il formalismo si generalizza immediatamente a un
caso dimensionale n :

$$\mathcal{L}[u(x_1, \dots, x_n), \frac{\partial u}{\partial x_i}, x_i]$$

$$\Rightarrow \int \mathcal{L} dx_1 \dots dx_n = 0$$

$$\Rightarrow \sum_{i=1}^n \frac{\partial}{\partial x_i} \frac{\partial L}{\partial (\frac{\partial u}{\partial x_i})} - \frac{\partial L}{\partial u} = 0$$

Ese. equazioni sull'onda in \mathbb{R}^3

$$\mathcal{L} = \frac{1}{2} \left[\left(\frac{\partial u}{\partial t} \right)^2 - (\vec{\nabla} u \cdot \vec{\nabla} u) \right] =$$

$$\frac{1}{2} \left[\left(\frac{\partial u}{\partial t} \right)^2 - \sum_{i=1}^3 \left(\frac{\partial u}{\partial x_i} \right)^2 \right]$$

potenzial $x_0 = t$ $x_1 = x, x_2 = y, x_3 = z$

$$\sum_{j=0}^3 \frac{\partial}{\partial x_j} \frac{\partial L}{\partial (\frac{\partial u}{\partial x_j})} - \frac{\partial L}{\partial u} = 0$$

$$\frac{\partial L}{\partial u} = 0$$

$$j=0 \Rightarrow \frac{\partial}{\partial t} \frac{\partial L}{\partial (\frac{\partial u}{\partial t})} = \frac{\partial^2}{\partial t^2} u$$

$$j=1 \Rightarrow \frac{\partial}{\partial x} \frac{\partial L}{\partial (\frac{\partial u}{\partial x})} = \frac{\partial}{\partial x} \left[\frac{\partial}{\partial (\frac{\partial u}{\partial x})} \left(\left(\frac{\partial u}{\partial x} \right)^2 + \left(\frac{\partial u}{\partial y} \right)^2 + \left(\frac{\partial u}{\partial z} \right)^2 \right) \right] \frac{1}{2}$$

$$= - \frac{\partial}{\partial x} \frac{\partial u}{\partial x} = - \frac{\partial^2 u}{\partial x^2}$$

$$\text{analog: } \frac{\partial}{\partial y} \frac{\partial L}{\partial (\frac{\partial u}{\partial y})} = - \frac{\partial^2 u}{\partial y^2} \quad \frac{\partial}{\partial z} \frac{\partial L}{\partial (\frac{\partial u}{\partial z})} = - \frac{\partial^2 u}{\partial z^2}$$

$$\Rightarrow \underbrace{\frac{\partial^2 u}{\partial t^2} - \frac{\partial^2 u}{\partial x^2} - \frac{\partial^2 u}{\partial y^2} - \frac{\partial^2 u}{\partial z^2}}_{-\Delta u} = 0$$

$$\boxed{\frac{\partial^2 u}{\partial t^2} - \Delta u = 0} \quad | \quad \text{Eq. aude in } \mathbb{R}^3$$